Stefan-Boltzmann's law of radiation with an amplifier

<table>
<thead>
<tr>
<th>Difficulty level</th>
<th>Group size</th>
<th>Preparation time</th>
<th>Execution time</th>
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<tr>
<td>hard</td>
<td>2</td>
<td>45+ minutes</td>
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</table>

Physics  Thermodynamics  Heat energy, thermal capacity


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General information

Application

Knowledge about the radiative behaviour of matter in dependence of temperature is widely used in fields such as astronomy or in the steel industry to check the temperature of the material, that is worked on.

Fig 1: Experimental set-up
According of Stefan-Boltzmann's law, the energy emitted by a black body per unit area and unit time is proportional to the power “four” of the absolute temperature of the body. Stefan-Boltzmann's law is also valid for a so-called “grey” body whose surface shows a wavelength-independent absorption-coefficient of less than one. In the experiment, the “grey” body is represented by the filament of an incandescent lamp whose energy emission is investigated as a function of the temperature.

The prior knowledge for this experiment is found in the Theory section.

The goal of this experiment is to check Stefan-Boltzmann's law of radiation.

1. To measure the resistance of the filament of the incandescent lamp at room temperature and to ascertain the filament's resistance $R_0$ at zero degrees centigrade.

2. To measure the energy flux density of the lamp at different heating voltages. The corresponding heating currents read off for each heating voltage and the corresponding filament resistance calculated. Anticipating a temperature-dependency of the second order of the filament-resistance, the temperature can be calculated from the measured resistances.
Theory (1/4)

If the energy flux density $L$ of a black body, e.g. energy emitted per unit area and unit time at temperature $T$ and wavelength $\lambda$ within the interval $d\lambda$, is designated by $dL(T, \lambda)/d\lambda$. Planck’s formula states:

$$\frac{dL(\lambda, T)}{d\lambda} = \frac{2\pi^2 \lambda^5}{c \hbar^3} e^{\frac{\hbar c}{\lambda kT}}$$

with $c = \text{velocity of light} = (3.00 \times 10^8 \text{[m/s]})$, $\hbar = \text{Planck’s constant} = (6.62 \times 10^{-34} \text{[J \cdot s]})$, $k = \text{Boltzmann’s constant} = (1.381 \times 10^{-23} \text{[J \cdot K^{-1}]})$

Integration of equation (1) over the total wavelength-range from $\lambda = 0$ to $\lambda = \infty$ gives the flux density $L(T)$ (Stefan-Boltzmann’s law).

$$L(T) = \frac{2\pi^2}{15} \cdot \frac{k^4}{\epsilon^3 h^3} \cdot T^4$$

Theory (2/4)

Respectively $L(T) = \sigma \cdot T^4$ with $\sigma = 5.67 \times 10^{-8}[\text{W \cdot m}^2 \cdot \text{K}^{-4}]$

The proportionality $L \sim T^4$ is also valid for a so-called “grey” body whose surface shows a wavelength-independent absorption-coefficient of less than one.

To prove the validity of Stefan-Boltzmann’s law, we measure the radiation emitted by the filament of an incandescent lamp which represents a “grey” body fairly well. For a fixed distance between filament and thermopile, the energy flux $\phi$ which hits the thermopile is proportional to $L(T)$.

$$\phi \sim L(T)$$

Because of the proportionality between $\phi$ and the thermoelectric e.m.f., $U_{\text{therm}}$ of the thermopile, we can also write:

$$U_{\text{therm}} \sim T^4$$
if the thermopile is at a temperature of zero degrees Kelvin. Since the thermopile is at room temperature $T_R$, it also radiates due to the $T^4$ law so that we have to write: $U_{therm} \sim (T^4 - T_R^4)$

Under the present circumstances, we can neglect $T_R^4$ against $T^4$ so that we should get a straight line with slope “4” when representing the function $U_{therm} = f(T)$ double logarithmically.

$$\lg U_{therm} = 4\lg T + \text{const.}$$

The absolute temperature $T = t + 273$ of the filament is calculated from the measured resistances $R(t)$ of the tungsten filament ($t =$ temperature in centigrade). For the tungsten filament resistance, we have the following temperature dependence:

$$R(t) = R_0(1 + \alpha t + \beta t^2) \quad \text{With } R_0 = \text{resistance at } 0^\circ C, \alpha = 4.82 \cdot 10^{-3} \cdot K^{-1}, \beta = 6.76 \cdot 10^{-7} \cdot K^{-2}$$

The resistance $R_0$ at $0^\circ C$ can be found by using the relation:

$$R_0 = \frac{R(t_R)}{1 + \alpha t_R + \beta t_R^2}$$

Solving $R(t)$ with respect to $t$ and using the relation $T = t + 273$ gives:

$$T = 273 + \frac{1}{2\beta} \left( \sqrt{\alpha^2 + 4\beta \left( \frac{R(t)}{R_0} - 1 \right)} - \alpha \right)$$

$R(t_R)$ and $R(t)$ are found by applying Ohm’s law, e.g. by voltage and current measurements across the filament.
## Equipment

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<th>Item No.</th>
<th>Quantity</th>
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<td>Optical bench expert l = 600 mm</td>
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<tr>
<td>2</td>
<td>Base for optical bench expert, adjustable</td>
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<td>Slide mount for optical bench expert, h = 80 mm</td>
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<td>Thermopile, Moll type</td>
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<td>Filament lamp 6V/5A, E14</td>
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<td>Connection box</td>
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Setup and Procedure

Setup

The experiment is started by setting up the circuit of Fig. 2 to measure the filament's resistance at room temperature. A resistor of 100 Ω is connected in series with the lamp to allow a fine adjustment of the current. For 100 mADC and 200 mADC the voltage drops across the filament are read and the resistance at room temperature is calculated. The current intensities are sufficiently small to neglect heating effects.

The experiment set-up of Fig. 1 is then built up. The 100 Ω resistor is no longer part of the circuit. The filament is now supplied by a variable AC-voltage source via an ammeter allowing measurement of alternating currents of up to 6 amperes. The voltmeter is branched across the filament and the alternating voltage is increased in steps of 1 volt up to a maximum of 8 V AC.

Figure 2: Circuit to measure the resistance of the filament at room temperature.
**Procedure**

**Remark:** the supply voltage of the incandescent lamp is 6 V AC. A voltage of up to 8 V AC can be applied if the period of supply is limited to a few minutes.

Initially, a voltage of 1 V AC is applied to the lamp and the Moll-thermopile, which is at a distance of 30 cm from the filament, is turned (slide-mount fixed) to the right and to the left until the thermoelectric e.m.f. shows a maximum. The axis of the cylindrical filament should be perpendicular to the optical bench axis. Since the thermoelectric e.m.f. is in the order of magnitude of a few millivolts, an amplifier has to be used for accurate readings. The factor of amplification will be 102 or 103 when using the voltmeter connected to the amplifier in the 10 V range. Before a reading of the thermoelectric e.m.f. is taken, a proper “zero”-adjustment has to be assured. This is done by taking the lamp together with its slide-mount away from the bench for a few minutes. The amplifier is used in the LOW DRIFT-mode (104 Ω) with a time constant of 1 s.

After the lamp has been put back onto the bench, the reading can be taken if the Moll-thermopile has reached its equilibrium. This takes about one minute. Care must be taken that no background radiation disturbs the measurement.

**Evaluation**
Task 1

Using the DC voltage output of the power supply unit, a direct current of 100 mA, respectively 200 mA, was supplied to the filament via a 100 $\Omega$ resistor. The corresponding voltage drops were found to be 16.5 mV and 33.0 mV. Doubling the current doubles the voltage drop. This shows that the temperature influence on the resistance is still negligibly small for the DC values chosen. We find in this case

$$R(t_R) = 0.165 \, \Omega$$

and hence $R_0 = 0.15 \, \Omega$

Small variations in $R_0$ only influence the slope $S$, which is to be found, in a negligible way.

Task 2

Increasing the AC heating voltage in steps of 1 V AC from 0 to 8 volts gave the following results:

The double logarithmic, graphical representation of the energy flux versus absolute temperature is shown in Fig. 3. The slope $S$ of the straight line is calculated, by regression, to be:

$$S = 4.19 \pm 0.265$$

The true value of $S$, which is 4, is found to be within the limits or error.

Figure 3: Thermoelectric e.m.f. of thermopile as a function of the filament's absolute temperature.